

A REVIEW ON MAGNETOHYDRODYNAMIC FLUID FLOWS

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ABSTRACT

In this review article a detailed study is presented on magnetohydrodynamic flows and their importance. Basic concept of MHD are discussed in detail. The results of various researchers are presented and addressed properly.

Keywords: MHD, Basic equations of MHD, Convective flows.

I. INTRODUCTION

The interaction of electromagnetic fields and fluids may be described from a scientific point of view, by the proper application of the principles of the special theory of relativity. The practical application of these principles, to actual physical phenomena of engineering, Astrophysics, Geo-physics, etc., is important in recent years. The study of this application to continuum has become known as magneto hydrodynamics or magneto fluid mechanics. The study of magneto hydrodynamics of viscous conducting fluids play a significant role during the recent times, owing to its practical interest and abundant applications in astro-physical and geo-physical phenomenon. Astrophysicists and geo-phycists realised soon after the advent of special relativity theory that electromagnetic fluid interactions were of great importance in stellar and planetary processes. The main impetus to the Engineering approach to electromagnetic fluid interaction studies has come from the concept of the MHD direct conversion generation, propulsion studies of radio propagation in the ionosphere, and controlled nuclear fusion. The study of flow problems of electrically conducting fluid particularly of ionized gases is currently receiving considerable interest. Such studies have made for many years in connection with astro-physical and geo-physical such as sun spot theory, motion of the interstellar gas etc. In recent years some Engineering problems need the studies of the flow of an electrically conducting fluid. Ionized gas is called plasma. Many names have been used in referring to the study of plasma phenomena. Harmann et al.[3] called it mercury dynamics since he worked with mercury. Astrophysicists called it cosmically electrodynamics and some called it magneto hydrodynamics. Physicists and electrical engineers commonly use the term plasma physics or plasma dynamics. The aerodynamicists have spoken of magneto hydrodynamics. The more general name magneto fluid mechanics was suggested by Vankarman et al. [12].

II. ELECTROMAGNETIC BODY FORCE

The basic concept on electromagnetism is Coulomb's law for the force on charged particles in an electrical field. Coulomb's law states that the force on charged particles is proportional to the product of the charge on the particle and the local electric field seen by the particle. In the rest frame of the particle and in RMKS units the law is stated as

$$\vec{F} = \rho_e \vec{E} \quad (1.1)$$

Where \vec{F} = Body force

\vec{E} = The electrical field And ρ_e = Space charge

When a charged particle moves in the presence of arbitrary electric and magnetic fields, the resulting force is called Lorentz force, which can be written as

$$\vec{F} = \rho_e (\vec{E} + \vec{q} \times \vec{B}) \quad (1.2)$$

Where \vec{q} = The velocity of the particle

And \vec{B} = The magnetic field seen by an observer in that frame.

The equation (1.2) is valid in any frame of reference, either in the rest frame (local frame of reference at rest with respect to the fluid) or the laboratory frame.

The electromagnetic body force results from the Coulomb force. The general quasi-state body force density for linear homogenous media can be written in the rest frame as

$$\bar{f}_e = \rho_e \bar{E} + \bar{J} \times \bar{B} \quad (1.3)$$

Where $\rho_e \bar{E}$ = The electric force

ρ_e = Space charge

$\bar{J} \times \bar{B}$ = The interaction force between the current and magnetic field.

Which becomes the predominant force in magneto fluid dynamics. This is called ponder motive force.

III. BASIC MHD EQUATIONS

The electromagnetic body force acts on the fluid and in turn the motion of the fluid in the presence of the electromagnetic field may generate an induced electromagnetic field and alter the field. Thus the velocity field in the fluid and electric and magnetic fields inside the fluid are disturbed due to the interaction of the fluid forces and the external electromagnetic force. We have to combine the basic concepts of fluid dynamics and electromagnetism to discuss the magneto fluid dynamics. The equations governing motion will be the Navier-Stokes equation of fluid and Maxwell's equations, which are usually coupled and must, solved simultaneously. The electromagnetic body force is usually non-conservative (rotational) and not derivable from a scalar potential function. Only under rare circumstances it may be approximately conservative and derivable from a scalar potential function. In such cases the electromagnetic body force alters simply the pressure, provided the boundary conditions on velocity and pressure remain the same. However this happens in exceptional cases and the force is rotational in general and does alter the flow. The Maxwell's equations in RMKS units are

$$\nabla \cdot \bar{D} = \rho_e \quad (\text{Coulombs law}) \quad (1.4)$$

$$\nabla \cdot \bar{B} = 0 \quad (\text{Absence of free magnetic poles}) \quad (1.5)$$

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} \quad (\text{Faraday's law}) \quad (1.6)$$

$$\nabla \times \bar{H} = \bar{J} + \frac{\partial \bar{D}}{\partial t} \quad (\text{Amper's law}) \quad (1.7)$$

And Ohm's law

$$\bar{J} = \sigma [\bar{E} + \bar{q} \times \bar{B}] + \rho_e \bar{q} \quad (1.8)$$

The current conservation equation is

$$\nabla \cdot \bar{J} + \frac{\partial \rho_e}{\partial t} = 0 \quad (1.9)$$

The continuity equation is

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \bar{q}) = 0 \quad (1.10)$$

The equation of motion (Navier-Stoke's equation) is

$$\frac{D\bar{q}}{Dt} = -\nabla p + \mu \nabla^2 \bar{q} + \bar{F} + \bar{f}_e \quad (1.11)$$

These equations hold in any frame of reference.

IV. MHD APPROXIMATIONS

In additions to the non-relativistic approximation, certain additional simplifications can be made for flow, which is quasi-steady (steady or low frequency oscillatory), and in which the electric field is of the order of magnitude of the induced quantity $\bar{q} \times \bar{B}$. Elsasser [62] has pointed out that the ratio of the displacement current to the conduction current in large class of phenomena is so small that displacement current can be neglected entirely. Thus both space charges and the electrostatic potential are omitted.

The following assumptions are made under MHD approximation.

(1) $\bar{q}^2 \ll C^2$. This allows the quantity $\sqrt{1 - \frac{\bar{q}^2}{C^2}}$ to be taken as unity and use the non-relativistic Newtonian form of equations of motion, where c is the velocity of the light.

(2) The electric field \bar{E} is of the same order of magnitude as the induced electric field $\bar{q} \times \bar{B}$. This assumption allows us to assume that the induced magnetic field is much smaller than the externally applied magnetic field.

(3) The problems of very high frequency are not considered, so that the displacement current $\frac{\partial \bar{D}}{\partial t}$ is neglected compared to \bar{J} , the conduction current.

Thus the equations of electromagnetic fields (1.6) and (1.7) become

$$\nabla \times \bar{E} = - \frac{\partial \bar{B}}{\partial t} \quad (1.12)$$

$$\nabla \times \bar{H} = \bar{J} \quad (1.13)$$

Equations (1.12). and (1.13) are called pre-Maxwell equations for electromagnetic fields.

(4) The electric energy is negligible compared to the magnetic energy. Since the displacement current and electric field energy are neglected, the main interaction is between the magnetic field and the fluid.

(5) The current density vector \bar{J} is approximately the same in any inertial frame. This means using Ohm's law we get

$$\bar{J} = \bar{J}' = \sigma (\bar{E} + \bar{q} \times \bar{B}) \quad (1.14)$$

(6) The magnetic force density represented by

$$\bar{f}_e = \rho_e \bar{E} + \bar{J} \times \bar{B} \quad (1.15)$$

The electric term can be shown to be negligible compared to magnetic term. For this reason the term $\rho_e \bar{E}$ is often neglected in MHD.

(7) In using equation $\nabla \cdot \bar{D} = \rho_e$ great care must be exercised. If the equation $\nabla \cdot \bar{J} + \frac{\partial \rho_e}{\partial t} = 0$ is used in metals (even in alternating current problem under the MHD approximation), the $\frac{\partial \rho_e}{\partial t}$ term may be dropped and

ρ_e need never be mentioned although it is identically zero.

Also $\nabla \cdot \bar{J} = 0$ is valid even for high frequency phenomena as long as the fluid velocity \bar{q} is very small compared to c .

V. EQUATIONS UNDER THE MHD APPROXIMATION

The basic equations under the MHD approximation take the following form.

$$\nabla \times \bar{H} = \bar{J} \quad (1.16)$$

$$\nabla \times \bar{E} = \frac{-\partial \bar{B}}{\partial t} \quad (1.17)$$

$$\nabla \cdot \bar{H} = 0 \quad (1.18)$$

$$\nabla \cdot \bar{J} = 0 \quad (1.19)$$

And the Ohm's law is

$$\bar{J} = \sigma (\bar{E} + \bar{q} \times \bar{B}) \quad (1.20)$$

The equation of conservation of mass is

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \bar{q}) = 0 \quad (1.21)$$

The equation of motion is

$$\rho \frac{D \bar{q}}{Dt} = -\nabla p + \mu \nabla^2 \bar{q} + \bar{J} \times \bar{B} - \rho \nabla \psi \quad (1.22)$$

Where ψ is the gravitational potential.

Magnetic diffusion equation is

$$\frac{\partial \bar{B}}{\partial t} = \eta \nabla^2 \bar{B} + \nabla \times (\bar{q} \times \bar{B}) \quad (1.23)$$

Where $\eta = \frac{1}{\sigma \mu_e}$ is magnetic diffusivity (or) magnetic viscosity.

The equation of energy is

$$\rho \frac{D(C_p T)}{Dt} = \phi - \rho (\nabla \cdot \bar{q}) + K \nabla^2 T + \frac{\bar{J}^2}{\sigma} \quad (1.22a)$$

In the case of free convection the momentum equation (1.22) reduces to

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = g_x - \beta(T - T_\infty) + \nu \frac{\partial^2 u}{\partial x^2} + \bar{J} \times \bar{B} \quad (1.24)$$

In the case of forced convection the momentum equation (1.22) becomes

$$\rho u \frac{\partial u}{\partial x} + p v \frac{\partial u}{\partial y} = -\frac{\partial p}{\partial x} + \frac{\partial}{\partial y} \left(\mu \frac{\partial u}{\partial y} \right) + \bar{J} \times \bar{B} \quad (1.25)$$

The energy equation (1.22a) reduces in both forced and free convection flows as

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \frac{\partial^2 T}{\partial y^2} + \frac{\bar{J}^2}{\sigma} \quad (1.26)$$

NON-DIMENSIONALISATION

We introduce the following dimensionless parameters

$$\begin{aligned} x^* &= \frac{x}{L}, & \nabla^* &= \nabla \bullet L, & t^* &= \frac{t}{t_0} = \frac{tU}{L}, \\ q^* &= \frac{\bar{q}}{U}, & \psi^* &= \frac{\psi}{gL}, & \bar{E}^* &= \frac{\bar{E}}{UB_0}, & \bar{H}^* &= \frac{\bar{H}}{H_0}, \\ \bar{J}^* &= \frac{\bar{J}}{\sigma\mu UH_0}, & p^* &= \frac{p}{\rho_0 U^2}, & \rho^* &= \frac{\rho}{\rho_0}, & \mu^* &= \frac{\mu}{\mu_0}, \end{aligned}$$

By using the above dimensionless parameters, the equations (1.16) to (1.23) after dropping the asterisks are

$$\nabla \times \bar{H} = R_\sigma \bar{J} \quad (1.27)$$

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} \quad (1.28)$$

$$\nabla \bullet \bar{H} = 0 \quad (1.29)$$

$$\nabla \bullet \bar{J} = 0 \quad (1.30)$$

$$\bar{J} = \sigma(\bar{E} + \bar{q} \times \bar{B}) \quad (1.31)$$

$$\frac{\partial \rho}{\partial t} + \nabla \bullet (\rho \bar{q}) = 0 \quad (1.32)$$

$$\rho \frac{D\bar{q}}{Dt} = -\nabla p + \frac{1}{R_e} \nabla^2 \bar{q} + \frac{M^2}{R_e} \bar{J} \times \bar{B} - \frac{\rho}{F_r} \nabla \psi \quad (1.33)$$

$$\frac{\partial \bar{B}}{\partial t} + -\frac{1}{R_\sigma} \nabla \times (\nabla \times \bar{B}) + \nabla \times (\bar{q} \times \bar{B}) \quad (1.34)$$

The dimensionless parameters appearing in the above equations are

$$(1) \quad \text{Re} = \text{Reynolds number} = \frac{UL}{\nu}, \text{ which is measure of the ratio of inertial to viscous forces.}$$

$$(2) \quad R_\sigma = \text{magnetic Reynolds number} = \frac{UL}{\eta} = \sigma\mu_e UL, \text{ which is a measure of the ratio of magnetic convection}$$

to magnetic diffusion. A small R_σ ($R_\sigma \ll 1$) indicates that the induced magnetic field is small compared to the total or the applied magnetic field.

$$(3) \quad F_r = \text{Froud number} = \frac{q_0^2}{gL} \text{ which is a measure of the ratio if inertial force to the gravitational force.}$$

$$(4) \quad M = \text{Hartmann number} = \sqrt{B_0^2 L^2 / \mu_r} = \sqrt{\sigma B_0^2 L^2 / \nu \rho_0} \text{ of the ratio of ponder motive force to the viscous forces. Here } \mu_f \text{ is the fluid viscosity.}$$

$$(5) \quad P_m = \text{magnetic Prandtl number} = \frac{\nu}{\eta} = R_\sigma / R_e, \text{ which is a measure of the ratio of vorticity diffusion to}$$

magnetic diffusion. This number is usually very small, so that the magnetic field diffuses much more easily than the vorticity.

MAGNATO HYDRODYNAMIC THERMAL CONVECTION

(a) MHD Convection equations

The following are the non-dimensional equations governing MHD convection flows:

$$\nabla \times \bar{H} = R_\sigma \bar{J} \quad (1.35)$$

$$\nabla \cdot \bar{H} = 0 \quad (1.36)$$

$$\nabla \times \bar{E} = -\frac{\partial \bar{B}}{\partial t} \quad (1.37)$$

$$\bar{J} = \sigma [\bar{E} + \bar{q} \times \bar{B}] \quad (1.38)$$

$$\nu M_0^2 p = \rho T \quad (1.39)$$

$$\nabla \cdot \bar{q} = 0 \quad (1.40)$$

We non-dimensional the momentum and energy equations (1.24), (1.25) and (1.26) using the dimensionless parameters defined in 6. The dimensionless form of the momentum and energy equations are

$$\frac{D\bar{q}}{Dt} = -\nabla p + \frac{1}{R_e} \nabla^2 \bar{q} + \frac{Gr}{R_e} \theta \hat{K} + \frac{M^2}{R_e} \bar{J} \times \bar{B} \quad (1.41)$$

$$\frac{D\theta}{Dt} = \frac{1}{P_e} \nabla^2 \theta + \frac{2E}{R_e} \varphi + \frac{M^2 E}{R_e} \bar{J}^2 \quad (1.42)$$

$$\text{And } \frac{\partial \bar{H}}{\partial t} = -\frac{1}{R_\sigma} \nabla \times (\nabla \times \bar{H}) + \nabla \times (\bar{q} \times \bar{H}) \quad (1.43)$$

Where the dimensionless heat function $\theta = \frac{T - T_0}{T_1 - T_0}$, E is Eckert

number, P_e is peck let number and Gr is Grashaff number.

(b) MHD Convective flows

The effect of viscous and the Joule dissipation on MHD free convection between two plates was first considered by Osterle et al.[7] and Poots et al. [8]. The former authors considered the case in which the loading conditions corresponds to a short circuit and the letter treated the open- circuit case and also the entrance flow.

Hughes et al. [5] considered the vertical channel to be connected through electrodes to an external circuit containing a finite load resistance RL and an emf source V_{app} and where the vertical plates are electrically conducting and of finite thickness. They assumed the electrical loading to be asymmetrical and employed an iteration scheme to solve the non-linear system of intergroddifferential equations governing this flow. Mahendra Mohan et al. [6] examined the free and forced convection effects in a rotating steady hydromagnetic viscous fluid between the two parallel plates maintaining the boundaries at constant temperature and taking the plates to be of finite conductivity. Prasada Rao et al. [9] investigated the influence of Hall currents on the free and forced convective flow of a viscous conducting fluid in a rotating channel maintained at constant temperature gradient along the channel walls under the influence of a transverse magnetic field. Rawat et al. [10] examined the free and forced convection effects on the flow of an electrically conducting viscous fluid between two parallel and electrically conducting porous plates subjected to suction and injection in the presence of applied magnetic field. Yang et al. [14] considered the free and forced convective MHD flow between two vertical plates. Harmann et al. [3] studied the problem of the steady viscous laminar flow of an electrically conduction liquid between parallel palates under a uniform transverse magnetic field. The problem is simple but gives insight into the MHD generator, pump, flow matter and bearings and it forms the basic method per treating all viscous flow devices.

Hughes et al. [4] considered the problem of flow through a rectangular channel bounded by walls, which are infinitely, conducting or perfectly insulating. The reason for considering such boundaries is the simplicity of the boundary conditions at such walls. The general problem of the rectangular channel flow has been discussed considering the various cases of insulating conducting walls. The cases of two parallel insulating walls and other two parallel conducting walls are the important MHD power generation. It is a general observation that an insulating electromagnetic force flattens the velocity profiles in a rectangular channel. Flows through and past porous media under the influence of magnetic field are of prime importance in ground water flow, petroleum, engineering and chemical engineering. Darcy law or Brinkman model governs flow in a porous medium. Flow in channel is described by Navier-Stokes equations. Chandrasekara et al.[1] examined the behaviour of the MHD flow in a squeeze film between two infinite strips where one of which has a porous bounding surface backed by rigid wall. Rudraiah and Chandrasekaran et al. [11] studied the MHD flow through a channel of varying gap. Vijayakumar Varma et al. [13] studied the steady MHD Poiseuille flow of a viscous conducting incompressible fluid in a channel bounded by porous media of which the lower wall is of high permeability and is bounded below by a rigid wall. Devika et al. [15] introduced the problem an analysis of first order homogeneous chemical reaction and heat source on MHD oscillatory flow of a visco-elastic fluid through a channel filled with saturated porous medium are reported. The present visco-elastic fluid model is working to suggest rheological liquids encountered in biotechnology (medical creams) and chemical engineering. This rheological model introduces additional terms into the momentum equation. It is assumed that the fluid has small electric conductivity and the electromagnetic force produced is very small. The dimensionless governing equations are solved analytically using regular perturbation method. The effects of various parameters on velocity, temperature and concentration fields are presented graphically and discussed. Raju et al. [16] discussed the problem of MHD free convective, dissipative boundary layer flow past vertical porous surface in the presence of thermal radiation, chemical reaction and constant suction, under the influence of uniform magnetic field which is applied normal to the surface is studied. The governing equations are solved analytically using a regular perturbation technique. The expressions for velocity, temperature and concentration fields are obtained. With the aid of these, the expressions for the coefficient of skin friction, the rate of heat transfer in the form of Nusselt number and the rate of mass transfer in the form of Sherwood number are derived. Finally the effects of various physical parameters of the flow quantities are studied with the help of graphs and tables. It is observed that the velocity and concentration increase during a generative reaction and decrease in a destructive reaction. The same observed to be true for the behavior of the fluid temperature. The presence of magnetic field and radiation diminishes the velocity and also the temperature. Ravikumaret al.[17] studied the problem of unsteady, two-dimensional, laminar, boundary-layer flow of a viscous, incompressible, electrically conducting and heat-absorbing Rivlin-Ericksen flow fluid along a semi-infinite vertical permeable moving plate in the presence of a uniform transverse magnetic field and thermal buoyancy effect is considered. The plate is assumed to move with a constant velocity in the direction of fluid flow while the free stream velocity is assumed to follow the exponentially increasing small perturbation law. Time-dependent wall suction is assumed to occur at the permeable surface. The dimensionless governing equations for this investigation are solved analytically using two-term harmonic and non-harmonic functions. The obtained analytical results reduced to previously published results on a special case of the problem. Numerical evaluation of the analytical results is performed and some graphical results for the velocity and temperature profiles within the boundary layer are presented. Skin-friction coefficient, Nusselt numbers are also discussed with the help of the graphs. Raju et al. [18] investigated the problem a steady MHD forced convective flow of a viscous fluid of finite depth in a saturated porous medium over a fixed horizontal channel with thermally insulated and impermeable bottom wall in the presence of viscous dissipation and joule heating. The governing equations are solved in the closed form and the exact solutions are obtained for velocity and temperature distributions when the temperatures on the fixed bottom and on the free surface are prescribed. The expressions for flow rate, mean velocity, temperature, mean temperature, mean mixed temperature in the flow region and the Nusselt number on the free surface have been obtained. The cases of large and small values of porosity coefficients have been obtained as limiting cases. Further, the cases of small depth (shallow fluid) and large depth (deep fluid) are also discussed. The results are presented and discussed with the help of graphs. Hari Priya et al. [19] discussed the combined effects of radiative heat transfer and a transverse magnetic field on steady

rotating flow of an electrically conducting optically thin fluid through a porous medium in a parallel plate channel taking hall current into account and non-uniform temperatures at the walls. The analytical solutions of velocity, temperature, shear stresses and rate of heat transfer are obtained from coupled nonlinear partial differential equations for the problem. The computational results are discussed quantitatively with the different variations in dimensionless parameters entering in the solution. Swetha et al. [20] found a perfect solution to the present natural convective flow problem of a vertical transfinite plate owing to the impulsive motion in the vicinity of first order chemical reaction, radiation absorption, radiation, Newtonian heating and species concentration in its plane is evolved by applying the method of Laplace transforms in closed form at the plate. Exact results for velocity, temperature, concentration fields are prevailed and expressions for heat and mass transfer rates are also found. The effects are analyzed for the respective invariables for both ammonia and water vapor. Misra et al. [21] considered the problem of oscillatory MHD flow of blood in a porous arteriole in presence of chemical reaction and an external magnetic field has been investigated. Heat and mass transfer during arterial blood flow are also studied. A mathematical model is developed and analyzed by using appropriate mathematical techniques. Expressions for the velocity profile, volumetric flow rate, wall shear stress and rates of heat and mass transfer have been obtained. Variations of the said quantities with different parameters are computed by using MATHEMATICA software. Ramaiah et al. [22] considered the combined influence of chemical reaction and radiation absorption effects on hydromagnetic free convective heat and mass transfer flow of viscous, incompressible, electrically conducting visco-elastic fluid through porous medium bounded by an oscillating porous plate in the presence of heat sources is studied. The expressions for velocity, temperature and concentration distribution, skin-friction, rate of heat and mass transfer coefficients at the plate are obtained using perturbation technique. The effect of various physical parameters occurring into the problem on velocity field is discussed with the help of graphs. Joseph et al. [23] considered the effect of variable suction on unsteady MHD oscillatory flow of Jeffrey fluid in a horizontal channel with heat and mass transfer has been studied. The temperature prescribed at plates is uniform and asymmetric. A perturbation method is employed to solve the momentum and energy equations. The effects of various dimensionless parameters on velocity and temperature profiles are considered and discussed in details through graphs. It is found that, the velocity u increases with increase in h_2 , α_1 , λ_1 , Gr , Gc , N , Re and Sc . The velocity also increases with decrease in h_1 , α_2 , Ha , and Kc . It is also observed that the temperature θ increases with increase in N and Pe . Increase in Schmidt number Sc and chemical parameter Kc respectively increase and decrease the species concentration or the concentration boundary layer thickness of the flow field. Maruf Hasan et al. [24] concentrated on the analysis of MHD free convection flow past an inclined stretching sheet. The viscous dissipation and radiation effects are assumed in the heat equation. Approximation solutions have been derived for velocity, temperature, concentration, Nusselt number, skin friction and Sherwood number using Nachtsheim-Swigert shooting iteration technique along with the six-order Runge-Kutta iteration scheme. Graphs are plotted to find out the characteristics of different physical parameters. The variations of physical parameters on skin friction coefficient, Nusselt number and Sherwood number are displayed via table. Falade et al. [25] investigated the effect of suction/injection on the unsteady oscillatory flow through a vertical channel with non-uniform wall temperature. The fluid is subjected to a transverse magnetic field and the velocity slip at the lower plate is taken into consideration. Exact solutions of the dimensionless equations governing the fluid flow are obtained and the effects of the flow parameters on temperature, velocity profiles, skin friction and rate of heat transfer are discussed and shown graphically. It is interesting to note that skin friction increases on both channel plates as injection increases on the heated plate. Padmavathi et al. [26] presented oscillatory plate temperature and dissipative effects on MHD free convection flow in a vertical. Padmavathi et al. [27] analyzed study of flow characteristics of conducting fluid embedded in porous media. Padmavathi et al. [28] proposed three-Dimensional Conducting Flow of Chemically Reactive Eyring-Powell Fluid over a Porous Stretching Sheet with Non-Uniform Heat Source / Sink and Soret Effects. Padmavathi et al. [29] demonstrated three-Dimensional Conducting Convective Flow of Williamson Fluid over a Porous Stretching Sheet With Chemical Reaction and Dufour Effects. Padmavathi et al. [30] examined numerical Study of an Electrically Conducting Three-Dimensional Casson Fluid Flow over a Porous Elastic Sheet with Non-Uniform Heat Source / Sink and Soret Effects. Padmavathi [31] presented oscillatory Plate Temperature and Unsteady Free Convection Flow of an Suction/ Injection Parameter on MHD in a Vertical Channel. Padmavathi

[32] analyzed three -Dimensional Hydromagnetic convective flow of chemically reactive Eyring- Powell Fluid With Non-Uniform Heat Source Sink. Padmavathi et. al [33] proposed three-Dimensional Hydromagnetic convective flow of Eyring-Powell fluid over a Bilinear Stretching Porous sheet with thermal Radiation and Chemical Reaction.

VI. CONCLUSIONS

Magnetohydrodynamic (MHD) fluid flow in different geometries relevant to human body parts is an interesting and important scientific area due to its applications in medical sciences. This article performs a comprehensive review on the applications of MHD and their numerical modelling in biological systems. Applications of MHD in medical sciences are classified into four categories. Applications of MHD in simple flow, peristaltic flow, pulsatile flow, and drug delivery are these categories. The numerical researches performed for these categories are reviewed and summarized separately. Finally, some conclusions and suggestions for future works based on the literature review are presented. The results indicated that during a surgery when it is necessary to drop blood flow or reduce tissue temperature, it may be achieved by using a magnetic field. Moreover, the review showed that the trapping is an important phenomenon in peristaltic flows that causes the formation of thrombus in blood and the movement of food bolus in gastrointestinal tract. This phenomenon may be disappeared by using a proper magnetic field. Finally, the concentration of particles that are delivered to the target region increases with an increase in the magnetic field intensity.

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